

Electric Charging of Dust Aggregates and its Effect on Dust Coagulation in Protoplanetary Disks

Satoshi Okuzumi (Kyoto U.)

Hidekazu Tanaka (Hokkaido U.)

Masa-aki Sakagami (Kyoto U.)



Ref.) Okuzumi (2009), ApJ, 698, 1122

Outline

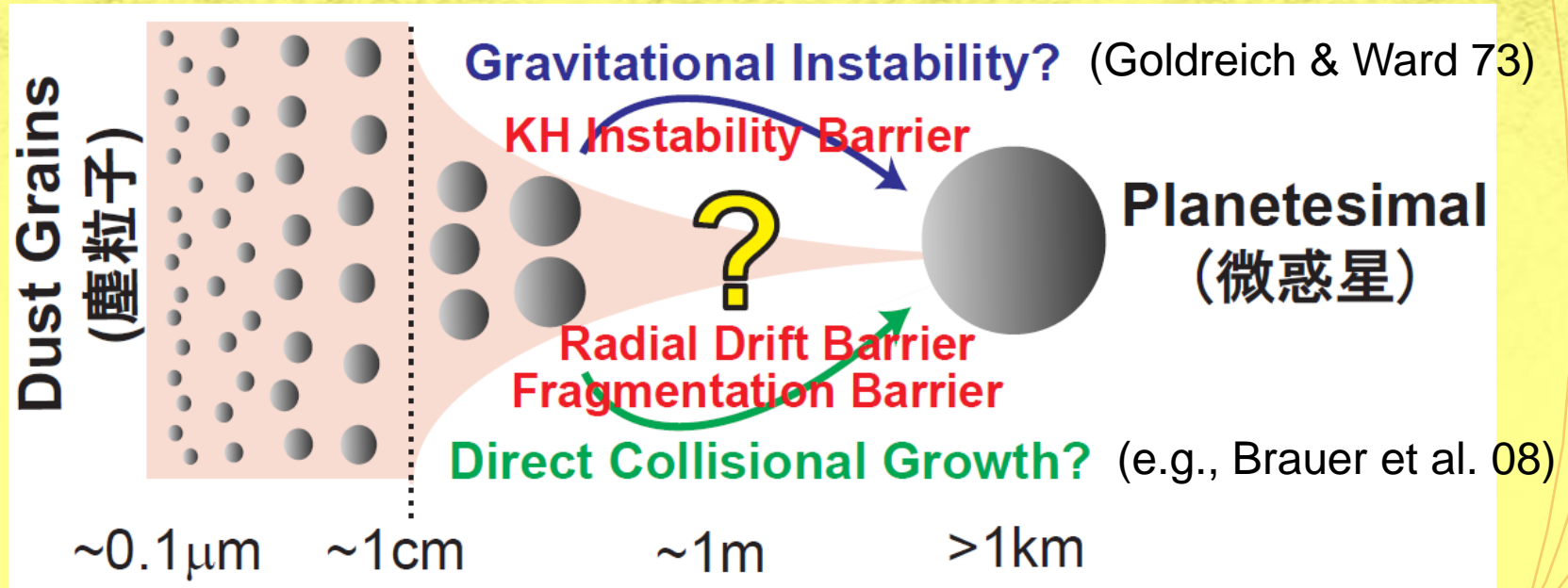
- ❑ Introduction: Dust charging in protoplanetary disks
 - ✓ Dust is an “MRI killer”

- ❑ Evolution of dust charge & gas ionization states
 - ← Analytical solutions

- ❑ “Electric barrier” against dust growth
 - ✓ Without turbulence, uniform (orderly) dust growth is likely to stall at its very early stage.

- ❑ Coagulation simulation with porosity evolution
 - Bimodal growth induced by the electric barrier

Dust Coagulation: The First Step toward Planet Formation



One way to go further:
More “realistic” modeling of dust aggregates

Porosity Evolution of Aggregates

“Classical” dust model:
compact sphere ($\rho = \text{const}$)

In reality, aggregates
can be very porous!

Low-velocity collision

→ Merger w/o restructuring

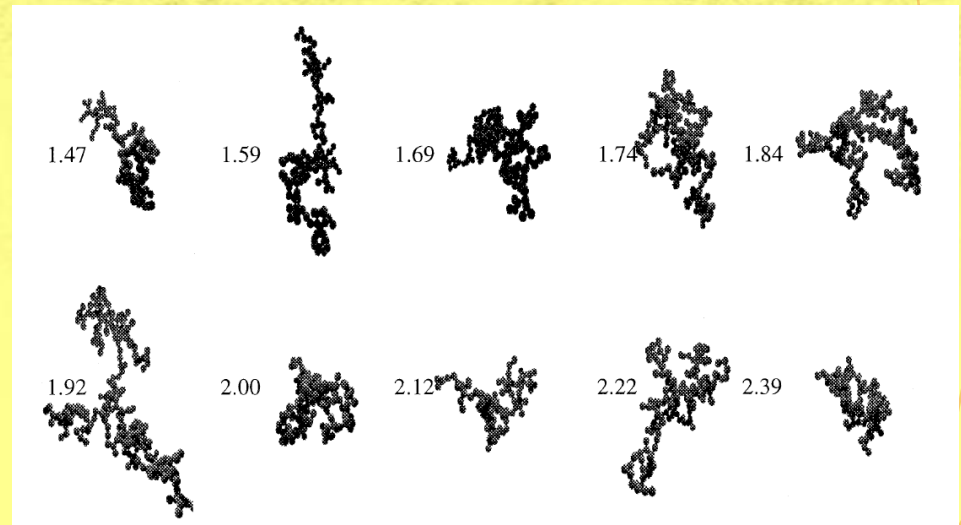
→ Fractal Structure: $D \approx 2$

$$a \approx a_0 N^{1/2}$$

✓ Fractal growth lasts

until porous size reaches 0.1-10cm (Blum 04, Suyama+ 08)

Experimental: Wurm & Blum 98; Blum+ 98
Statistical: Ormel+ 07; Zsom & Dullemond 08
N-body: Kempf+ 99; Wada+ 07,08; Suyama+08



An N-body result of thermal coagulation
(Kempf et al. 1999)

High porosity → Large surface area

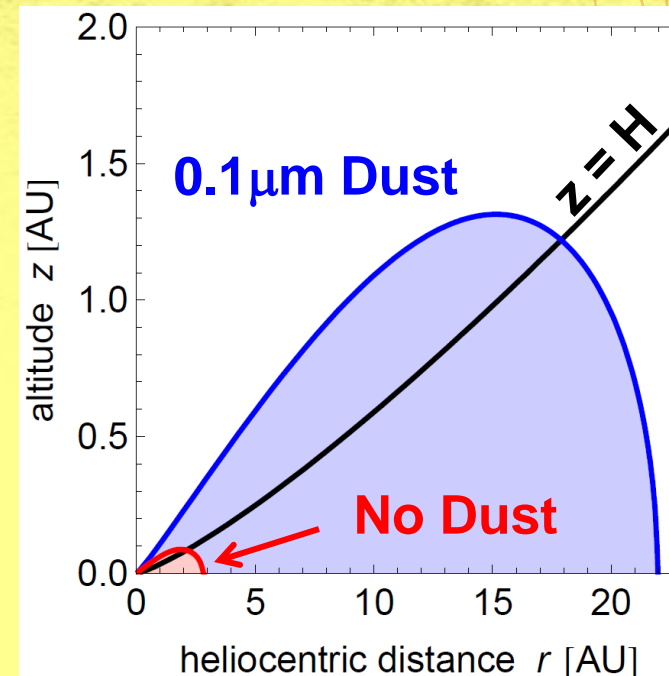
→ Strong coupling to gas, **Efficient ion & electron capturing**

Another Key: Dust Charging

- MRI provides disk turbulence which may cause:
 - ✓ Fragmentation of large aggregates
 - ✓ Stirring-up of small aggregates/fragments
- Small grains/aggregates greatly reduces x_e by negative charging
 - Wide “dead zone” (Sano+ 00)

Dust = “MRI Killer”

Location of the Dead Zone
MMSM+CR+XR+RA



Dust Growth & Settling \Leftrightarrow MRI-driven Turbulence
Interact via Dust Charging

Question:

What's the role of the electrostatic interaction of aggregates in their collisional growth?

- 1: How the dust charge state (& gas ionization state) evolve with dust growth?
- 2: How the dust charging feed back to the growth?

Step 0: Dust Charging in Neutral Plasmas

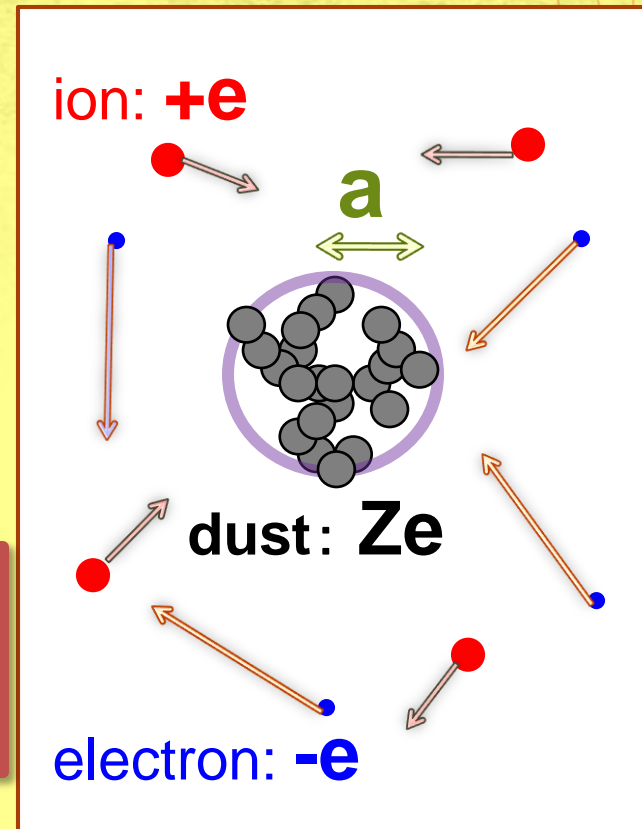
Thermal velocity: **Electron** \gg **Ion**



Negatively Charged

Charge Equilibrium Condition:
$$-\frac{Ze^2}{a} \sim k_B T$$

➔
$$Z \sim -\frac{ak_B T}{e^2} \sim \left(\frac{a}{0.1\mu\text{m}}\right) \left(\frac{T}{100\text{K}}\right)$$



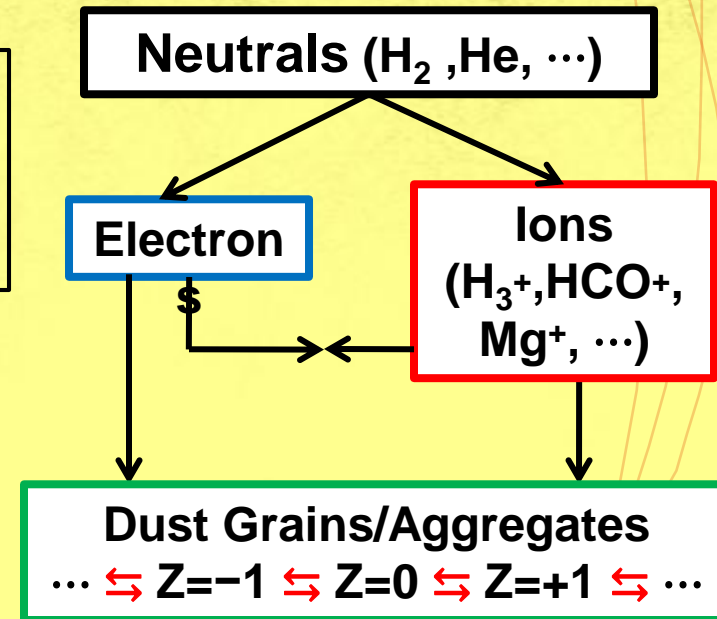
- An exact solution exists for spherical bodies (Spitzer 1941)
- Roughly holds for fluffy aggregates (Matthews & Hyde 2008)

Step 1: Charge State in Weakly Ionized Gas

- ★ Dust charge state : $n_d(Z)$
 - ★ Gas ionization state : n_i & n_e
- Must be computed **consistently**

- ❑ Previous Work (Sano+ 00, Ilgner & Nelson 06)
 - Solved charge transfer **numerically**
 - **Heavy to compute**

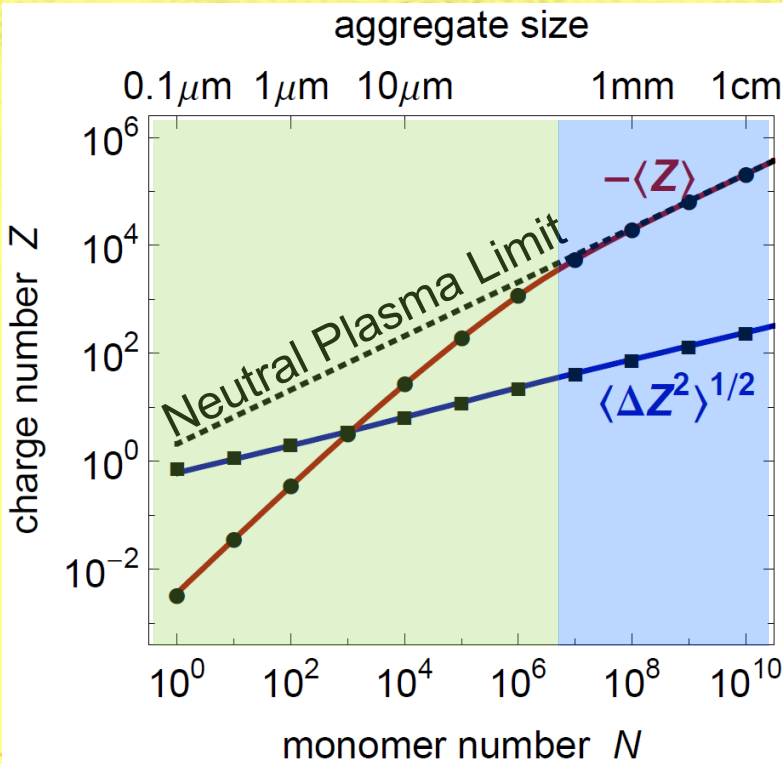
- ❑ This Work: solves charge transfer **as analytically as possible**
 - ✓ Single eq. for a “master” parameter
 - ✓ Arbitrary size & porosity distribution
 - ✓ Agrees with numerical solutions very well (see next slide)
 - ✓ Enables coupled simulation of dust growth/settling & MRI turbulence



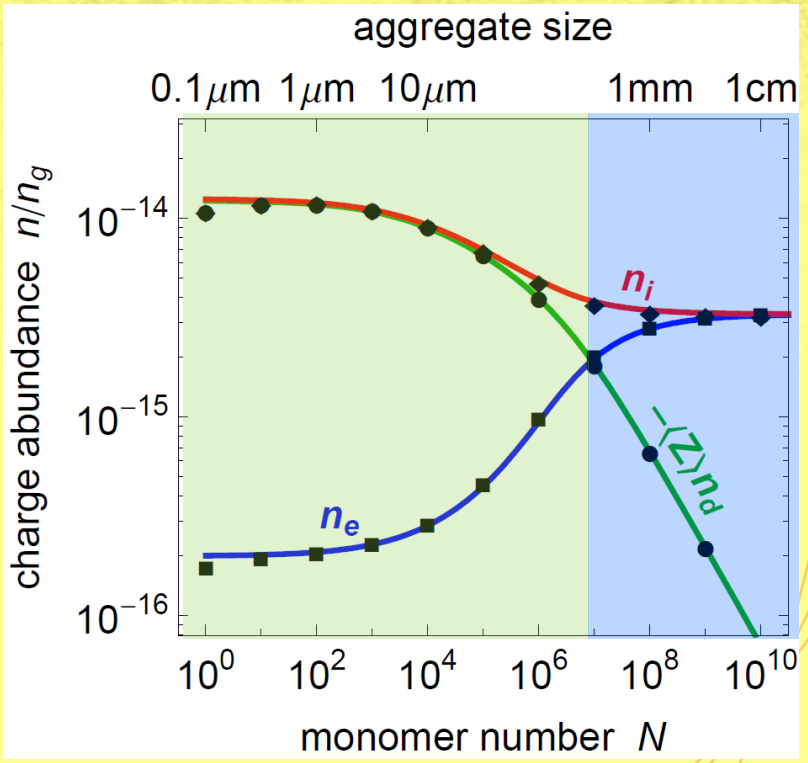
Co-Evolution of Dust Charge & Gas Ionization States

Example: $r = 5\text{AU}$, $z = \text{H}$, **fractal growth ($D=2$)** : $a \approx a_0 N^{1/2}$

Evolution of Dust Charge State



Evolution of Gas Ionization State



Small Dust → **Ion-Dust Plasma (IDP)** : $n_i \approx -\langle Z \rangle n_d$,

$$\frac{\langle Z \rangle e^2}{a} \ll k_B T$$

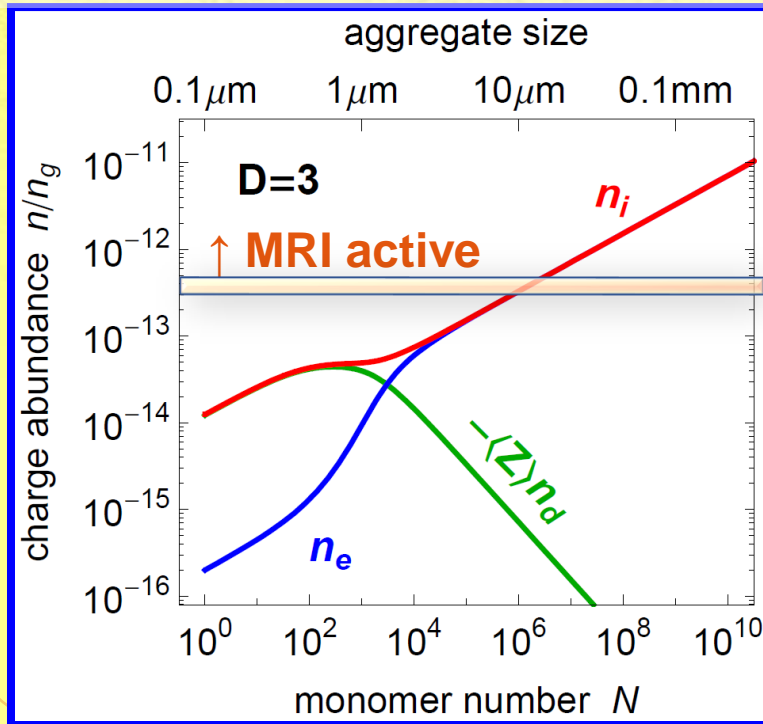
Large Dust → **Ion-Electron Plasma (IEP)** : $n_i \approx n_e$,

$$\frac{\langle Z \rangle e^2}{a} \sim k_B T$$

Fractal Dust Maintains a Wide Dead Zone

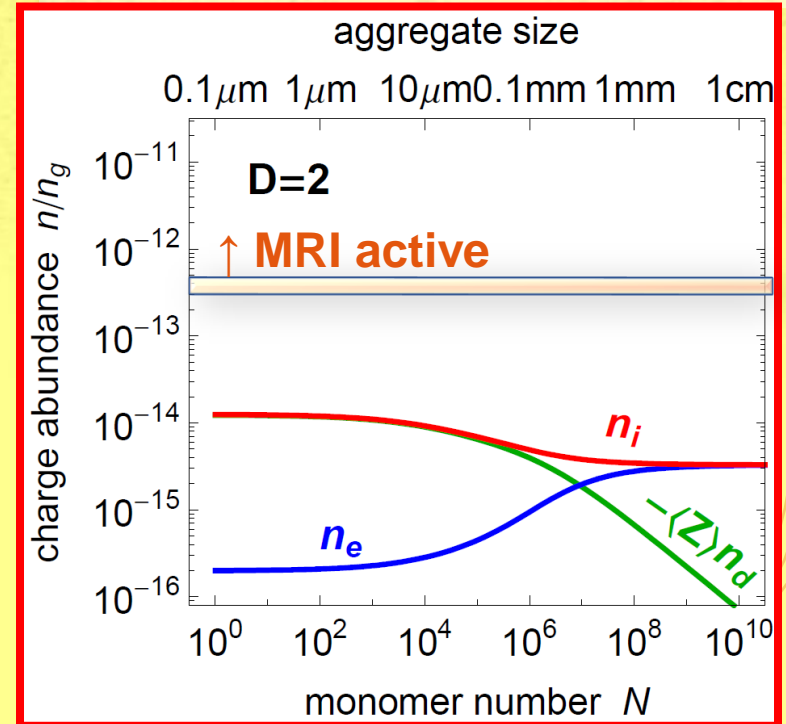
$D \leq 2 \rightarrow$ Total surface area σn_d conserved $\rightarrow n_e$ kept low

Compact ($D=3$) Dust Model



MRI active for size $a > 10 \mu\text{m}$

Fractal ($D=2$) Dust Model



MRI inactive for all sizes!!

\rightarrow Dead zone is maintained until aggregates are compressed ($a \sim \text{cm}$)

Step 2: A Simple Estimation of the Effect of Electrostatic Repulsion on Dust Growth

$$E_{\text{kin}} = \frac{1}{2} \mu v_{\text{rel}}^2$$

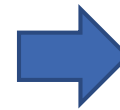
Relative
Kinetic Energy

$$E_{\text{el}} = \frac{Z Z' e^2}{a + a'}$$

“Electric Energy”
before Contact

Collisional
Cross Section:

$$\sigma_{\text{eff}} = \pi(a + a')^2 \left(1 - \frac{E_{\text{el}}}{E_{\text{kin}}}\right)$$



Growth Possible if

$$E_{\text{kin}} > E_{\text{el}}$$

- ✓ Dust Motion: Brownian Motion + Sedimentation + Turbulence
- ✓ Growth Mode: **Monodisperse(=uniform) & Fractal (D=2)**
 ← valid for early evolutionary stages ($a < 1\text{cm}$)
- ✓ Disk Model: MMSN + Ionization sources
 (cosmic ray + X-ray + radionuclides)

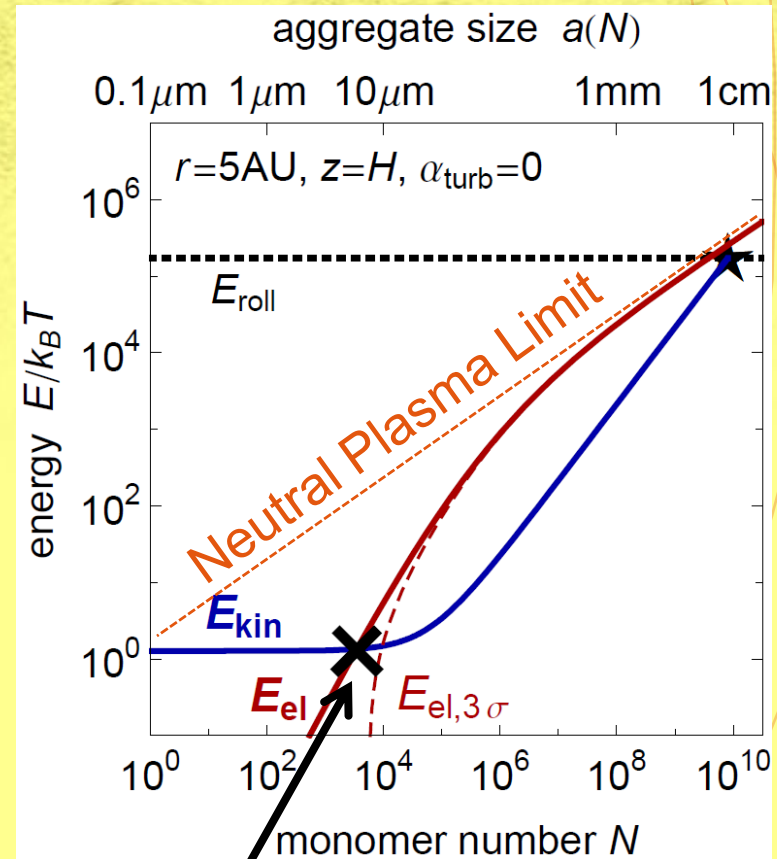
Kinetic Energy vs Electrostatic Energy

Example: $r = 5\text{AU}$, $z = H$, no turbulence

$$E_{\text{el}} : Z_1 = Z_2 = \langle Z \rangle$$

$$E_{\text{el},3\sigma} : Z_1 = \langle Z \rangle, Z_2 = \langle Z \rangle + \underline{3\langle \Delta Z^2 \rangle^{1/2}}$$

Both E_{el} & $E_{\text{el},3\sigma}$ quickly converge to the neutral plasma limit, exceeding E_{kin} much before collisional compaction (★) begins!

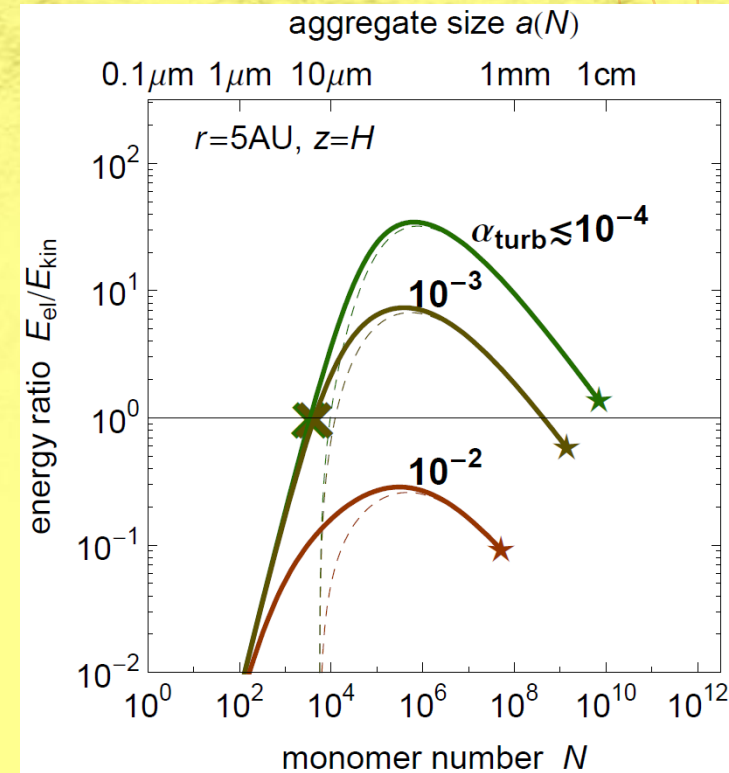


Uniform growth stalls at size as small as $10\mu\text{m}$!

Effect of Turbulence

- Disk turbulence enhances the collisional velocity Δv .
(see, e.g., Ormel & Cuzzi 07)
- The electric barrier ($E_{\text{el}} > E_{\text{kin}}$) is removed for $\alpha \gtrsim 10^{-2}$.

→ If MRI turbulence is present, dust is likely to overcome the electric barrier.



But...

Disk Turbulence: a Friend or Foe?

At later stages, such strong turbulence is very likely to cause **collisional fragmentation** of aggregates!!

□ $\Delta v \sim \alpha^{1/2} c_s \sim 100 \text{ m/s}$ for $t_{\text{fric}} \sim T_k$ & $\alpha \sim 10^{-2}$
(Ormel & Cuzzi 07)

□ Collision of aggregates results in fragmentation for $\Delta v > 30\text{-}60 \text{ m/s}$
(Wada+ 08,09(submitted))

Laminar → Electric barrier at early growth stages
Turbulent → **Fragmentation barrier at late stages**

Summary up to here

- ❑ Evolution of dust charge & gas ionization states with dust growth is investigated analytically.
 - ✓ Fractal growth (valid for size $< \text{cm}$) tends to maintain a large dead zone.

- ❑ “Electric barrier” against dust growth:
 - ✓ Without turbulence, uniform (orderly) dust growth is likely to stall at its very early stage.
 - ✓ With turbulence, dust can overcome the barrier, but then suffers from collisional fragmentation.

A Possible Path to Evolve: Bimodal Growth

Small “field” aggregates (mass m)
 Large “test” aggregates (mass M)

$$M \gg m$$

$$n(M) \ll n(m)$$

➔ Test aggregates can collide with “frozen” field aggregates if

$$(\Delta v)_{Mm} > (\Delta v)_{mm}$$

Brownian-motion dom. ➔ $(\Delta v)_{Mm} < (\Delta v)_{mm}$

Sedimentation-dom. ➔ $(\Delta v)_{Mm} \gtrsim (\Delta v)_{mm}$

A small fraction of massive aggregates are allowed to continue growing even if the growth of the others has “frozen out.”

Step 3: Simulation of Porous Dust Coagulation

- Extended Smoluchowski Equations:
 - Evolve mass DF $n(M)$ & mass-volume rel. $V(M)$ consistently
- Porosity change ← Empirical formula from N-body simulations

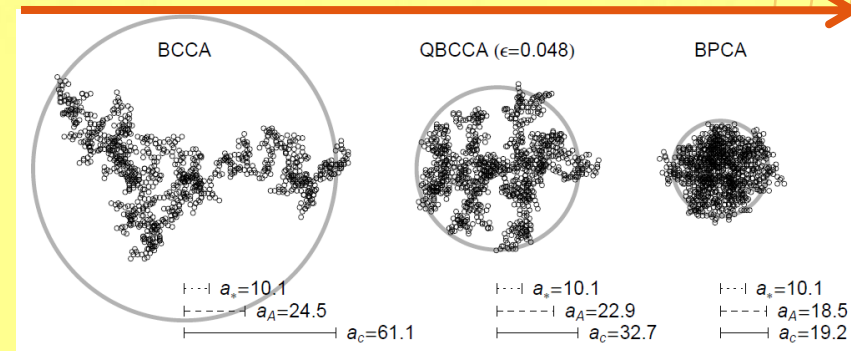
(Okuzumi, Tanaka, & Sakagami, coming soon!)

$$\partial_t n_M = \int_0^{M/2} \frac{K_{M-M', V_{M-M'}}^{M', V_{M'}}}{n_M} n_{M'} n_{M-M'} dM'$$

$$- n_M \int_0^\infty \frac{K_{M', V_{M'}}^{M, V_M}}{n_{M'}} n_{M'} dM'$$

$$\partial_t V_M = \int_0^{M/2} [V_{1+2}(V_{M'}, V_{M-M'}) - V_M] \times \frac{K_{M-M', V_{M-M'}}^{M', V_{M'}}}{n_{M'}} n_{M'} n_{M-M'} dM'$$

Decreasing V_m/V_M



Small Size Ratio



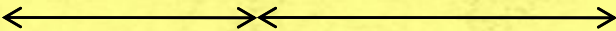
Small Porosity Increase

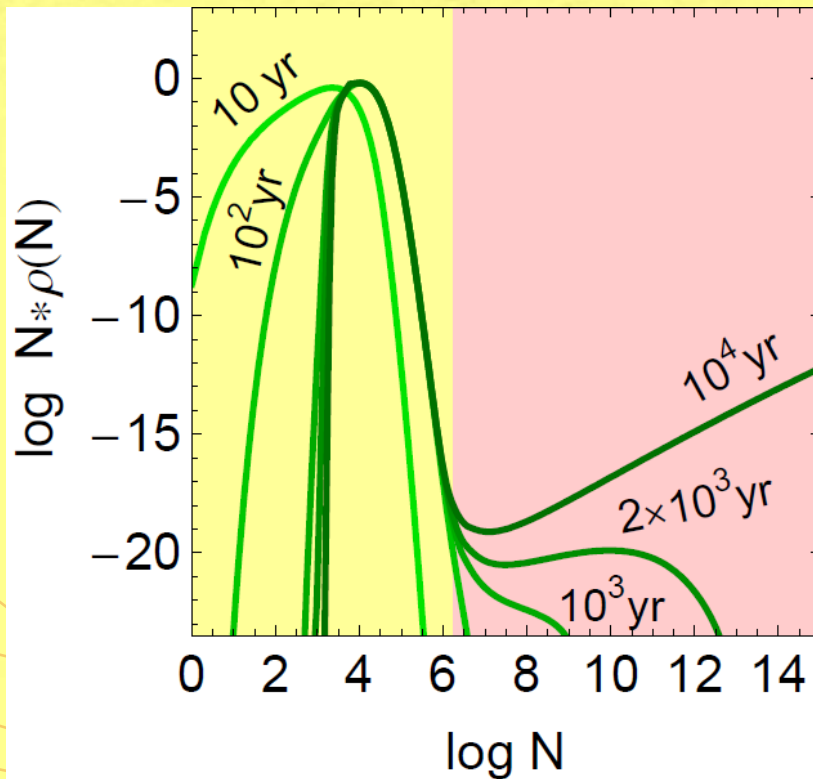
$V_{1+2}(V_m, V_M)$: volume of collisional outcome

Result

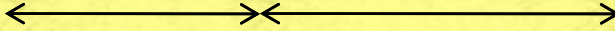
$r = 5\text{AU}$, $z = H$, $\alpha = 0$, Ionization Rate = $0.1 \times (\text{X-ray Ionization})$

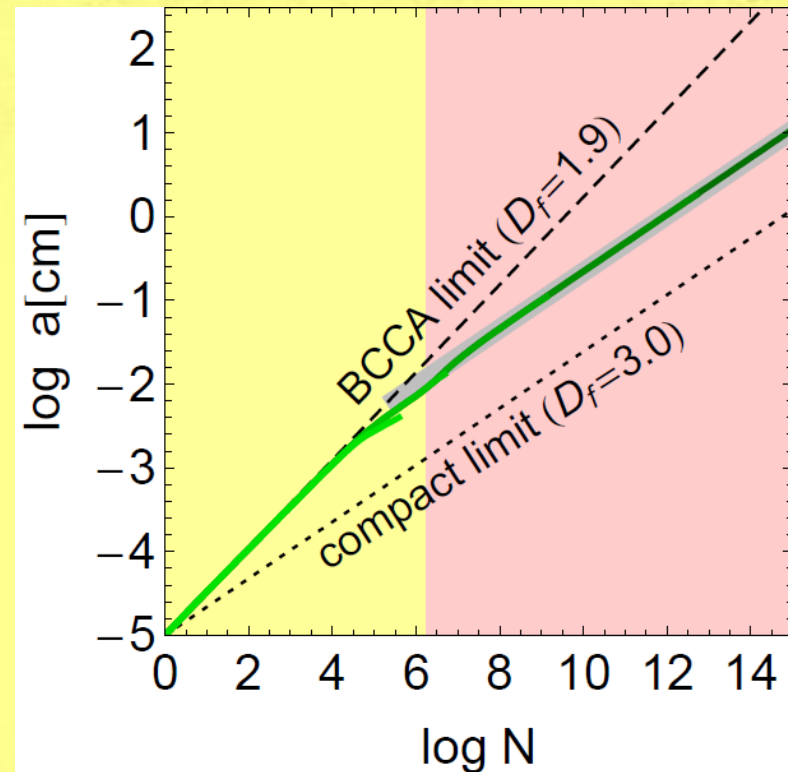
Mass Spectrum

Brownian dom. Sedimentation dom.




Mass-Radius Relation

Brownian dom. Sedimentation dom.




Summary

- ❑ Evolution of dust charge & gas ionization states with dust growth is investigated analytically.
 - ✓ Fractal growth (valid for size < cm) tends to maintain a large dead zone.

- ❑ “Electric barrier” against dust growth
 - ✓ Fractal growth is likely to stall in the absence of strong turbulence

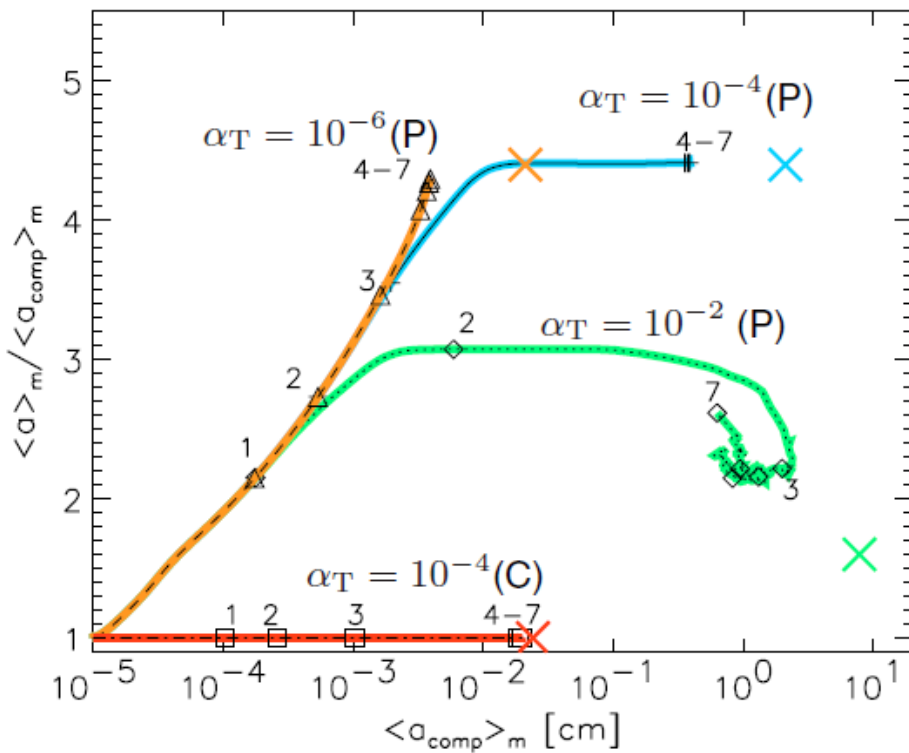
- ❑ Bimodal growth induced by the electric barrier is confirmed by a coagulation simulation including porosity evolution.

Code Check: Comparison with a Full-2D Method

VAM reproduces full-2D results in surprisingly good accuracy

Full-2D Monte Carlo Method (Ormel+ 07)

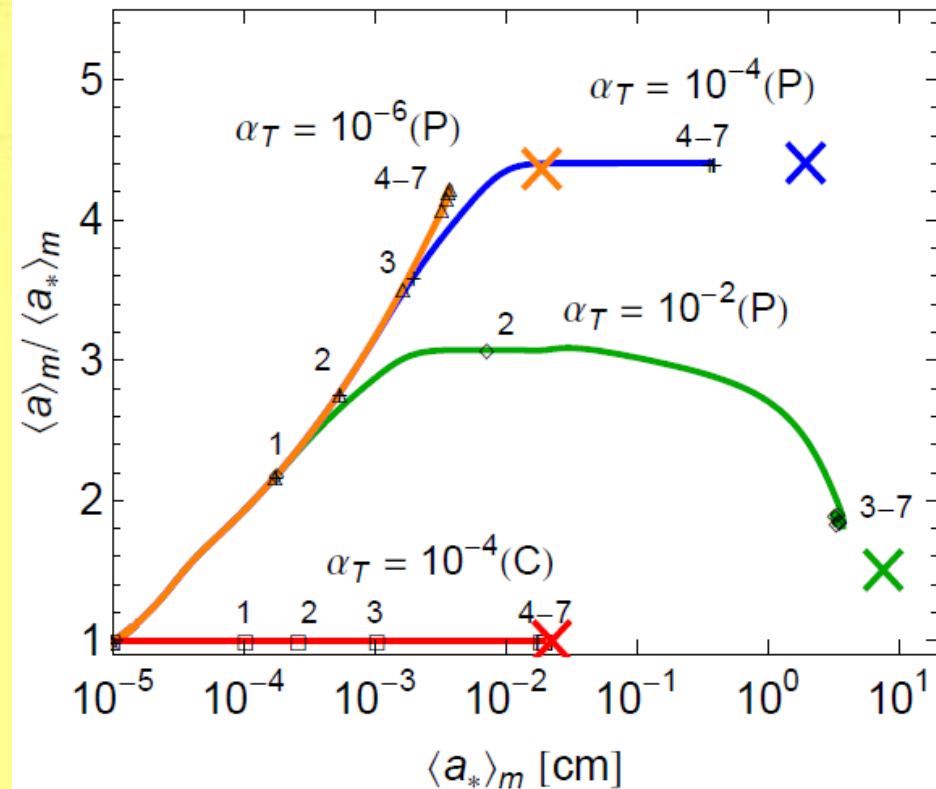
Evolution curves 1 AU models



CPU Time: 2 Hours

Volume-Averaging Method

Evolution curves 1 AU models



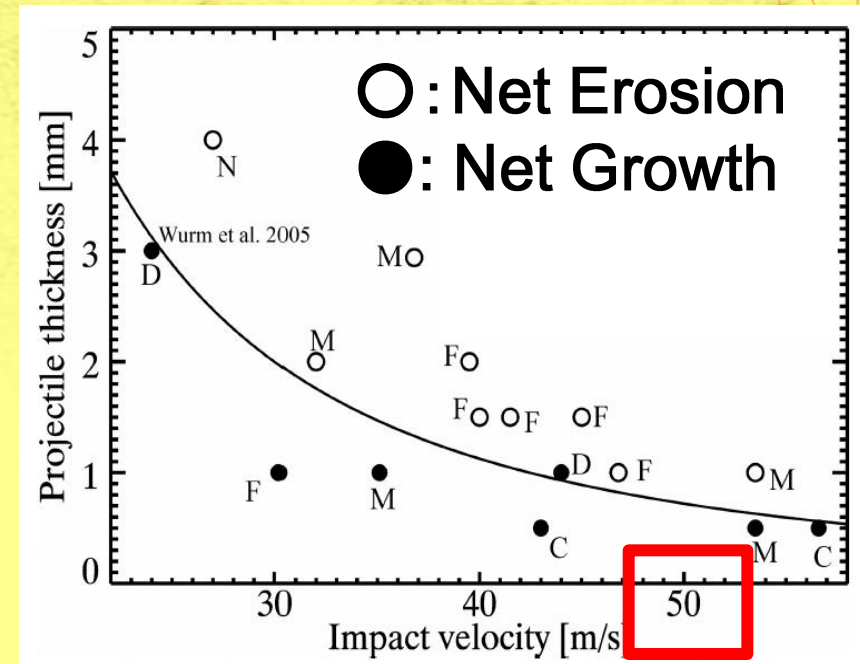
CPU Time: 1 Min

Is Bimodal Growth Favorable for Planetesimal Formation?

Recent Lab Experiment (Teiser & Wurm 08)

Collision of large (1-10cm) target and small (<mm) projectile

→ Target can achieve net growth even if $v_{\text{imp}} > 50 \text{ cm/s}$!!



Teiser & Wurm (2008)

Bimodal growth might overcome the fragmentation barrier!

Porous Dust Can Overcome the Radial Drift Barrier

$r = 5\text{AU}, z = 0$

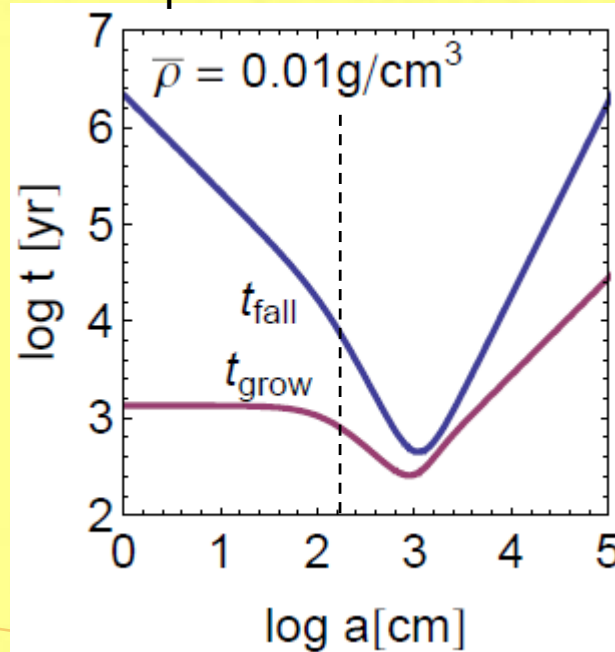
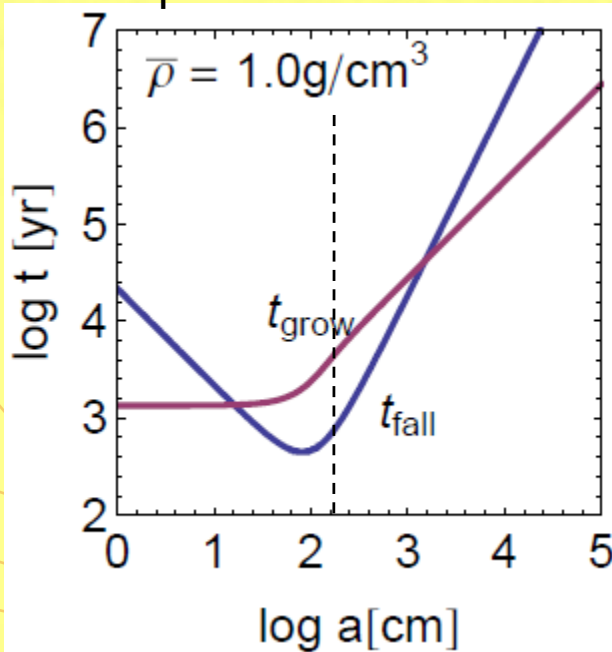
$$t_{\text{stop}} \propto \frac{ma}{\sigma} \propto \bar{\rho} a^2$$

Compact Dust Model
(Porosity = 0%)

Porous Dust Model
(Porosity = 99%)

Epstein Stokes

Epstein Stokes



$$t_{\text{stop}} \sim \Omega_K^{-1} \Rightarrow a_c \propto \bar{\rho}^{-1/2}$$

$$\left\{ \begin{array}{l} \sigma \propto \bar{\rho}^{-1} \\ m \propto \bar{\rho} a_c^3 \propto \bar{\rho}^{-1/2} \end{array} \right.$$

$$\frac{\dot{m}}{m} \propto \frac{\sigma}{m} \propto \bar{\rho}^{-1/2}$$